



Synthesis of CoFe_2O_4 superparamagnetic nanoparticles using a rapid thermal processing furnace with halogen lamps

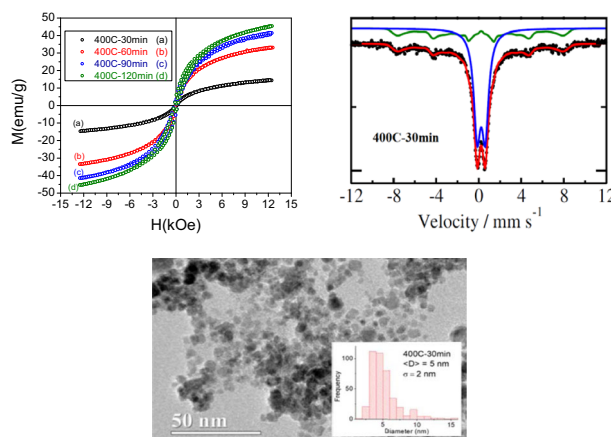
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Abstract

This work deals with the synthesis of superparamagnetic cobalt ferrite (CoFe_2O_4) nanoparticles, via a sol-gel method that uses gelatin as an organic precursor and a rapid thermal processing furnace with halogen lamps as a heat source. TEM, HRTEM, XRD, VSM and Mössbauer spectroscopy measurements were performed, at room temperature, to characterize the samples. Nanoparticles with an average size of 5–10 nm and microstrain of the order of 10^{-4} – 10^{-3} were obtained. Magnetic hardening was also observed with increased nanoparticle size, in addition to an increase in the anisotropy constant, which ranged from 1.7 to $4.7 \cdot 10^6$ erg/cm³. The saturation magnetization decreased with decreasing size due to the formation of a magnetically dead layer on the surface (thickness = 0.6–0.7 nm). The Mössbauer spectra showed an increase in the superparamagnetic phase for samples with smaller nanoparticles and less dispersion.

Graphical Abstract



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Keywords Nanoparticles · Cobalt ferrite · Superparamagnetic · Sol–gel preparation · Halogen lamps furnace

Highlights

- The protein sol–gel method using gelatin is a fast and low-cost process for the production of nanoparticulate powders.
- The use of RTP furnaces in calcination helps to control the energy supplied to the sample.
- The advantages of using halogen lamps in RTP furnaces are low cost, uniform heating, and good temperature control.
- In order to have a superparamagnetic material, it is necessary to produce nanoparticles with a size below a critical size.

Abbreviations

RTP	rapid thermal processing
RTP-HL	halogen lamp RTP systems
Sch	Scherrer's formula
WHP	Williamson–Hall plot
SSP	size–strain plot
XRD	X-ray diffraction
TEM	transmission electron microscopy
HRTEM	high-resolution transmission electron microscopy
VSM	vibrating-sample magnetometry

1 Introduction

The production and control of the size and morphology of nanoparticles are some of the most researched subjects in the scientific community [1–4]. Producing nanoparticles with an increasingly smaller size with a specific shape, using the least amount of energy, are the main objectives [5–7]. Reducing the size of the nanoparticles may allow the appearance of a new magnetic property or reinforce an existing property [4, 8]. One of the ways to achieve this is by controlling the heat supply during synthesis, as heat favors the growth of nanoparticles [5–8]. The use of RTP (rapid thermal processing) systems is highly recommended in these cases as they have high heating rates, thus reducing the calcination time, in addition to having high cooling rates [5, 6, 8–12]. Halogen lamp RTP systems (RTP-HL) can achieve a heating rate of up to hundreds of K/s, while conventional systems can only reach a few tens of K/min [5]. The advantages of using halogen lamps are: low cost, uniform heating, and good temperature control [5]. RTP systems are also used to control the crystallization of amorphous phases, control phase transitions, reduce strains in the crystalline structure, produce crystallographic texture [5], change coercivity [8, 9, 11], and increase magnetic permeability [12].

Zeng et al. [9] used the RTP technique with halogen lamps, whose heating rate was 200 °C/min, to show that FePt nanoparticles produced in the RTP furnace had greater coercivity than those produced in a conventional furnace. Corbiere et al. [7] produced vanadium nitride nanoparticles using a halogen lamp furnace and whose heating rate was

50 K/min, obtaining nanoparticles smaller than 10 nm. Chu et al. [6] produced nanocomposites with NdFeB and realized that the samples produced in an RTP system, when compared to samples produced in a conventional furnace, showed a reduction in grain size, greater morphological uniformity in addition to an increase in coercivity.

It is known that cobalt ferrite has cubic magnetocrystalline anisotropy, several axes of easy magnetization, reasonable saturation magnetization, positive anisotropy constant, and good mechanical and chemical stability. Due to these characteristics, it is well suited for the production of permanent and isotropic magnets, magnetic recording, magnetic fluids, guided drug delivery, medical diagnoses, hyperthermia treatment, etc. For this to be possible it is necessary to produce stable particles with a single magnetic domain. It is estimated that to occur the formation of a single domain it is necessary that the particles have a size between 7–10 nm. In form bulk, cobalt ferrite is a hard magnetic material with a Curie temperature around 790 K [1, 3, 4]. Grigorova et al. [1] showed in his work that cobalt ferrite nanoparticles have saturation magnetization ranging from 13–70 emu/g, remanence ranging from 1–23 emu/g, and coercivity ranging from 0.07 to 1.75 kOe. Kumar et al. [3] showed in their work that the saturation magnetization ranged from 2 to 50 emu/g, the remanence ranged from 0.5 to 25 emu/g and coercivity ranged from 0.5 to 1.4 kOe. Rajendran et al. [4] showed in his work that the saturation magnetization ranged from 9 to 49 emu/g, the remanence ranged from 0.55 to 7 emu/g and zero coercivity.

2 Experimental

The samples were prepared by the protein sol-gel method that uses gelatin as an organic precursor [13]. Initially 4.000 g of gelatin was dissolved in 10 mL of distilled water at room temperature. Then 8.000 g of $\text{Fe}(\text{NO}_3)_3 \cdot 9\text{H}_2\text{O}$ was also dissolved in 10 mL of distilled water, at room temperature. The two solutions were then mixed in beaker I. The same sequence was performed with 2.8820 g of $\text{Co}(\text{NO}_3)_2 \cdot 6\text{H}_2\text{O}$ and 1.4419 g of gelatin and placed in a beaker II. The solutions of beakers I and II were then mixed and subjected to magnetic stirring for 1 h at room temperature.

The final solution was put to dry in a furnace at 100 °C for 48 h. The xerogel formed was macerated and calcined at 400 °C for 30, 60, 90, and 120 min in the RTP furnace with halogen lamps (heating rate: 100 °C/min). Due to a cooling system, the furnace returned to room temperature in ~1 h. The sample calcined was washed with hydrogen peroxide to eliminate the remaining organic matter. Then it was centrifuged and dried in a furnace at 100 °C for 48 h. Finally, X-ray diffraction (XRD), transmission electron microscopy (TEM), high-resolution transmission electron microscopy (HRTEM), vibrating-sample magnetometry (VSM), and Mössbauer spectroscopy measurements were made.

Cobalt ferrite (bulk) has a cubic unit cell, with space group Fd-3m (227) and cell parameters $a = b = c = 0.83919$ nm (ICDD 00-022-1086). Its structure can be inverse spinel or mixed spinel, the latter occurring more frequently [3, 14, 15].

The XRD measurements were performed at room temperature in a PanAlytical diffractometer, model X'Pert Pro (MPD), operated at 40 kV and 40 mA, with a cobalt target ($\lambda_{K\alpha 1} = 0.1789$ nm) and hybrid monochromator. The diffraction patterns were obtained in the range 2θ : 10°–100° with 0.013° step and a speed of 100 s per step. The crystalline phases were identified using the X'Pert HighScore Plus 3.0d software. TEM and HRTEM measurements were performed at room temperature in a transmission electron microscope model Jeol (JEM-2100), equipped with EDS, Thermo Scientific. The VSM measurements were performed at room temperature using a vibrating sample magnetometer whose magnetic field ranged from –12 kOe to +12 kOe. The Mössbauer spectra were measured at room temperature in a SEECO Spectrometer, Model W302, with a radioactive source diffused in a rhodium matrix and were adjusted with NORMOS software, using the hyperfine field distribution method.

3 Results and discussion

From the XRD measurements (Fig. 1a), it is observed that the cubic cobalt ferrite phase was formed in all samples. The average size D of the nanoparticles was determined by Scherrer's Formula (Sch) $D = k\lambda/(\beta\cos\theta)$ [16], by a linear fit of the Williamson–Hall Plot (WHP) $\beta\cos\theta/\lambda = k/D + 4\epsilon\sin\theta/\lambda$ [17] and by the linear fit of the Size–Strain Plot (SSP) $(d\beta\cos\theta/\lambda)^2 = (k/D)^2d^2 + 2\pi\langle\epsilon\rangle^2$ [18]. The microstrain ϵ and its root mean square value, $\langle\epsilon\rangle$, were calculated from linear fits to the WHP and SSP, respectively. As the nanoparticles were approximately spherical in shape, the value chosen for the form factor was $k = 1$ [16]. The value of d (d -spacing) is calculated by Bragg's Law. It can be seen from Table 1 that the size of the nanoparticles increased with increasing calcination time. This is because

the growth of nanoparticles and coalescence are stimulated by the supply of energy. It is also observed that the microstrain oscillated, despite the fact that it was expected to decrease with the increase of the calcination time, since the crystalline structure approaches the ideal with the energy supply [19].

Single-domain particles exhibit superparamagnetic behavior below a certain critical size D_C [15, 19]. For a system of spherical, single-domain and uniaxial nanoparticles, of CoFe_2O_4 , with low size dispersion, the value of D_C is in the range 8–10 nm [4, 15, 20]. In Fig. 1b and c we have the TEM and HRTEM images of samples 400C-30 min and 400C-120 min, respectively, where it is possible to observe the approximately spherical shape of the nanoparticles, with the sample 400C-30 min having a smaller average and less dispersion (smaller standard deviation σ). Through HRTEM images it was possible to visualize crystalline planes and to identify them from the d -spacing: 0.29 nm (220) and 0.48 nm (111). Figure 2a and Table 1 show that saturation magnetization (M_S), maximum field magnetization (M_{max}), remanant magnetization (M_R) and coercivity (H_C) increased with the increase in the size of the nanoparticles and this increased the number of blocked particles. Consequently, we have a gain in the ferrimagnetic phase for samples with a larger nanoparticle size. By making a fit to the magnetization curve, by the saturation approximation law $M = M_S[1 - 0.0762K^2/(M_S^2H)]$ [21], it was possible to determine the values of M_S and K (anisotropy constant) for each sample. The anisotropy constant, which represents an energy density, increased with increasing average size, since larger nanoparticles have more atoms and thus store more energy due to magnetization. The K value fell in the range $1.7\text{--}4.7 \cdot 10^6$ erg/cm³, which is in line with the results found by Machado et al. [19], Babić-Stojić et al. [22], and Mooney et al. [15]. According to Vázquez-Vázquez et al. [21], magnetic nanoparticles do not have uniform magnetization, but their magnetization decreases from the center to the surface, and on the surface there is a layer, of thickness t , with zero magnetization, called the dead layer. It is possible to determine the value of this thickness t using the dead layer model $M_s = M_s(\text{bulk})[1 - 6t/\langle D \rangle]$ [21], where M_s (bulk) is the saturation magnetization of a bulk sample and $\langle D \rangle$ is the average size measured by TEM images. Figure 2b shows that $t = 0.6$ nm and $M_s(\text{bulk}) = 89.11$ emu/g, which is in accordance with the results obtained by the TEM measurements in Fig. 1, and those obtained by Vázquez-Vázquez et al. [21] and by Mooney et al. [15]. For samples with the lowest average size values, this layer is responsible for significant reductions in saturation magnetization and maximum field magnetization (see Fig. 2 and Table 1).

Figure 3 shows the ZFC (zero-field cooled) and FC (field cooled) curves of the 400C-30 min sample measured

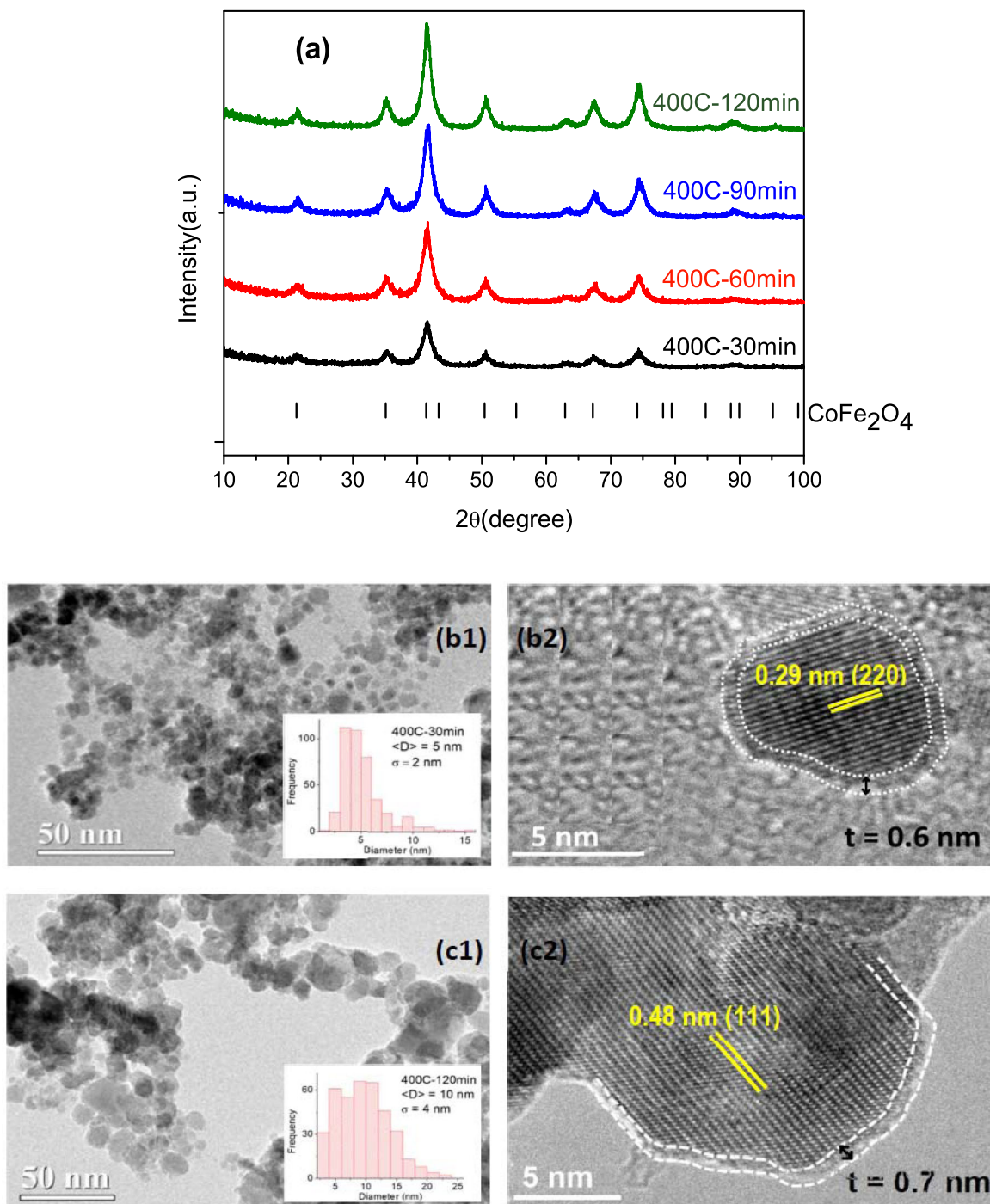


Fig. 1 a XRD of the four samples (ICDD 00-022-1086). b and c TEM, HRTEM and histograms of samples 400C-30 min and 400C-120 min, respectively

Table 1 Average size, microstrain and magnetic measurements

400 °C	Average size (nm)				Microstrain		Magnetic measurements				
							emu/g			kOe	-10^6 erg/cm^3
	TEM	Sch	WHP	SSP	WHP	SSP	M_S	M_{\max}	M_R	H_C	K
30 min	5	6	7	7	$2 \cdot 10^{-3}$	$3 \cdot 10^{-3}$	18	15	0.04	0.04	1.7
60 min	6	6	6	6	$3 \cdot 10^{-4}$	$2 \cdot 10^{-3}$	39	33	1.20	0.06	2.9

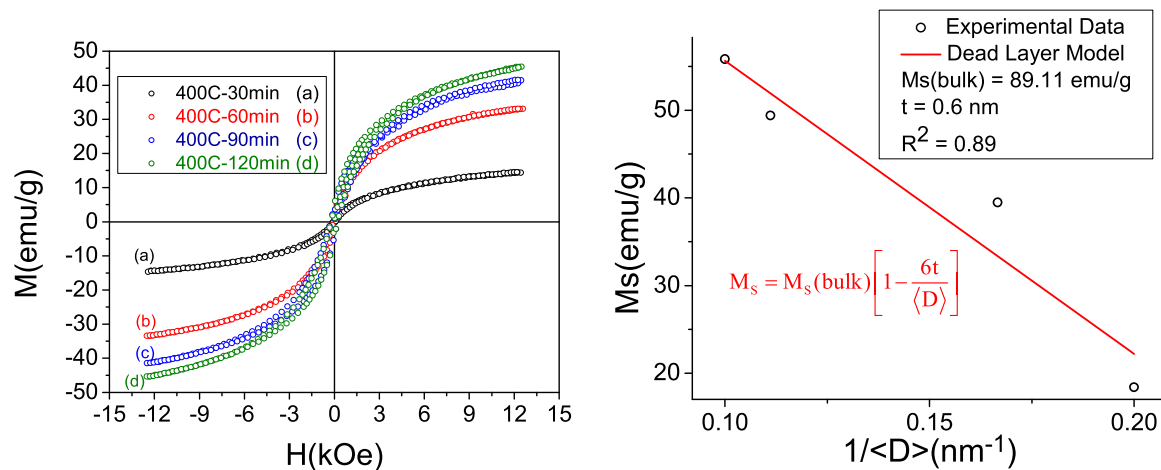


Fig. 2 **a** M versus H measurements of the four samples and **b** Linear Fit of the Dead Layer Model

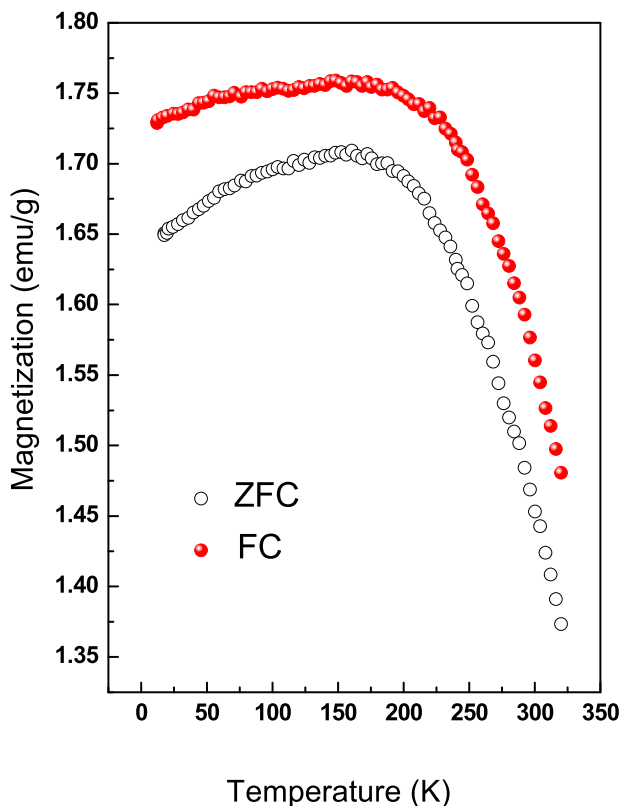


Fig. 3 ZFC–FC curves of sample 400C-30 min under 100 Oe magnetic field

in the temperature range between 10 and 330 K under the application of a 100 Oe magnetic field. As the temperature increases, the magnetization referring to the ZFC curve also increases until reaching a maximum magnetization value of 1.71 emu/g around 165 K (blocking temperature). However, even after the transition from the blocked to the superparamagnetic state, the ZFC–FC curves do not overlap to reach a temperature of magnetic irreversibility.

Therefore, these results indicate that we have, at room temperature, a system composed of strongly interacting superparamagnetic nanoparticles.

Figure 4 shows the Mössbauer spectra, measured at room temperature, for all samples. The 400C-120 min sample showed a distribution of hyperfine fields (B_{hf}) centered at 48T, characteristic of the α -magnetite phase [23]. However, it is possible to observe in the spectra of samples 400C-90 min, 400C-60 min, and 400C-30 min a doublet with quadrupole splitting values of 0.77, 0.76, and 0.90 mm/s, characteristic of the superparamagnetic phase, induced by the effect of reducing the size of the nanoparticles [23]. In addition, the ferrimagnetic contributions of the samples are not completely suppressed, since, in all spectra, a distribution of B_{hf} is identified, showing the superposition of the ferrimagnetic and superparamagnetic phases. The analysis of the spectral area relationship between the doublet and the sextet shows that the superparamagnetic phase present in samples 400C-90 min, 400C-60 min, and 400C-30 min corresponds to 15%, 44%, and 68%, respectively. Such behavior is a reflection of the reduction in size, which induces the superparamagnetic phase. Such results are in agreement with the TEM, XRD, and VSM analyses.

4 Conclusion

It was concluded that the sol–gel process, together with the RTP-HL furnace, was efficient in producing superparamagnetic nanoparticles of cobalt ferrite. It was observed that the increase in the calcination time increased the size of the nanoparticles and caused magnetic hardening, but the microstrain underwent unexpected changes. The increase in size of the nanoparticles also caused an increase in the saturation magnetization

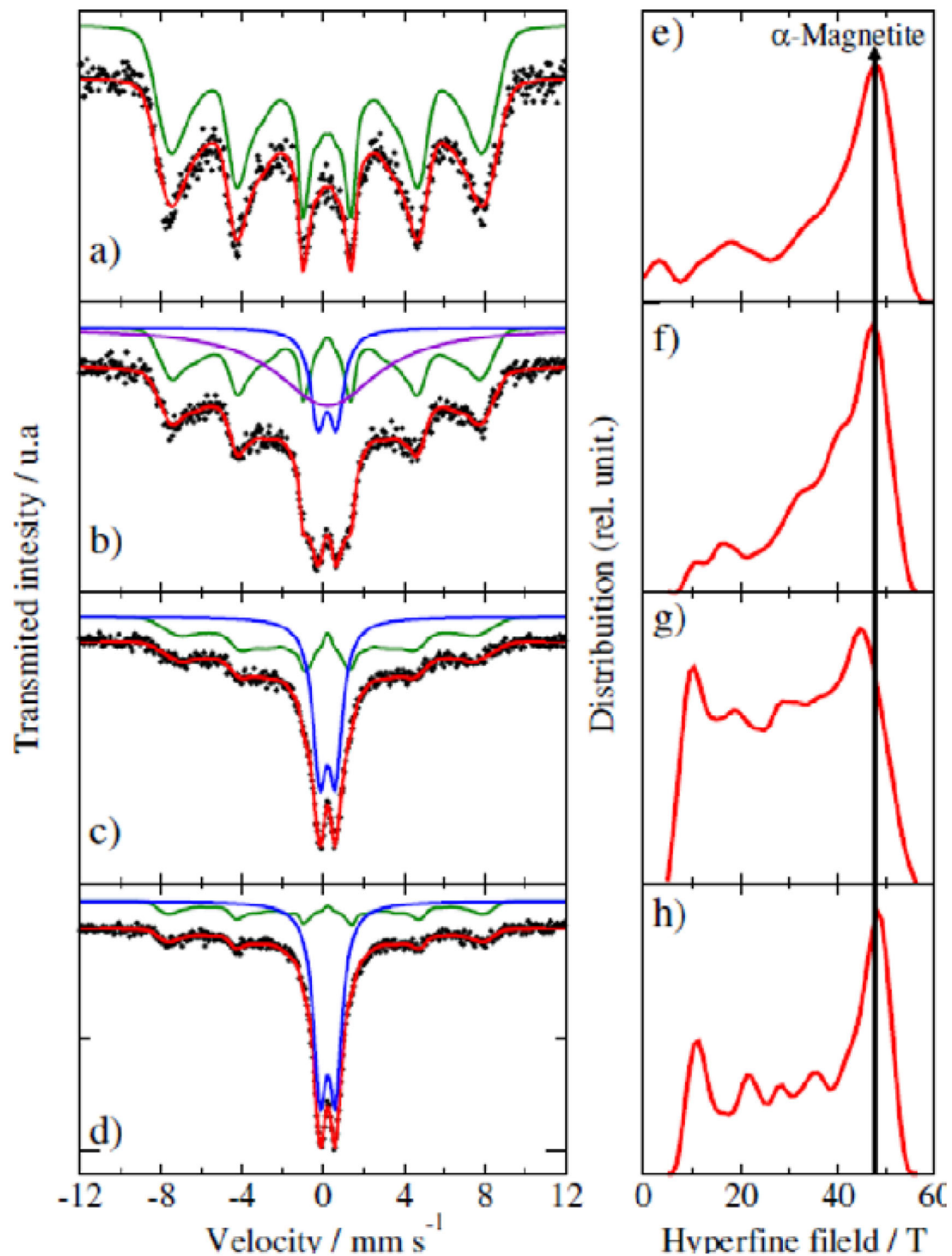


Fig. 4 Mössbauer spectra of the samples and distribution of hyperfine fields. **a–e** 400C-120 min, **b–f** 400C-90 min, **c–g** 400C-60 min, and **d–h** 400C-30 min

and in the anisotropy constant. The thickness of the dead layer that covered the nanoparticles was estimated to be 0.6 nm. Magnetization measurements indicated a blocking temperature close to 165 K. The results of the Mössbauer

spectra showed that the samples with the smaller average size (5–6 nm) have a higher fraction of the superparamagnetic phase and, consequently, lower coercivity values.

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